

Optimization approximation with separable variables for the one-way wave operator

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The thin slab propagator

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$$\begin{aligned}\partial_k \hat{p} + s\rho \hat{v}_k &= \hat{f}_k \\ s\kappa \hat{p} + \partial_r \hat{v}_r &= \hat{q}\end{aligned}$$

where

$$\hat{p}(x_m, s) = \mathcal{L}[p(x_m, t)]$$

κ compressibility modulus

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- ▶ Reduced system
- ▶ The coupled system of one-way equations
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The thin slab propagator

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The thin slab propagator

- ▶ The thin slab propagator
- ▶ Reduced system to handle horizontal velocity of the particle

$$\hat{v}_\mu = -\rho^{-1} s^{-1} (\partial_\mu \hat{p} - \hat{f}_\mu)$$

- ▶ The coupled system of one-way equations
- ▶ The propagator

The thin slab propagator

- ▶ The thin slab propagator
- ▶ Reduced system
from which we obtain the matrix differential equation

$$(\partial_3 \delta_{I,J} + s \hat{A}_{I,J}) \hat{F}_J = \hat{N}_I$$

$$\hat{F}_1 = \hat{\rho} \quad , \quad \hat{F}_2 = \hat{v}_3$$

$$\hat{A}_{1,1} = \hat{A}_{2,2} = 0$$

$$\hat{A}_{1,2} = \rho$$

$$\hat{A}_{2,1} = -D_\nu(\rho^{-1} D_\nu) + \kappa$$

$$\hat{N}_1 = \hat{f}_3 \quad , \quad \hat{N}_2 = D_\nu(\rho^{-1} \hat{f}_\nu) + \hat{q}$$

- ▶ The coupled system of one-way equations
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- ▶ The coupled system of one-way equations

$$(\partial_3 \delta_{I,M} + s \hat{\Lambda}_{I,M}) \hat{W}_M = -(\hat{L}^{-1})_{I,M} (\partial_3 \hat{L}_{M,K}) \hat{W}_K + (\hat{L}^{-1})_{I,M} \hat{N}_M$$

- ▶ The propagator

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For the coupled system, we can get the operator

$$(\hat{G}\hat{u})(x_\mu, x_3) = \int_{\zeta \in \mathcal{R}} \int_{x'_\nu \in \mathcal{R}} \hat{\Gamma}(x_\mu, x_3; x'_\nu, \zeta) \hat{u}(x'_\nu, \zeta) dx'_1 dx'_2 dx'_3 d\zeta$$

with the initial conditions

$$(\partial_3 + s\hat{\Gamma})\hat{U} = 0 \quad , \quad \text{for } x_3 > \zeta \quad , \quad \hat{U}(x_\mu, \zeta; \zeta) = \hat{u}(x_\mu, \zeta)$$

we get

$$(\hat{G}\hat{u})(x_\mu, x_3) = \int_{\zeta=-\infty}^{x_3} \hat{U}(x_\mu, x_3; \zeta) d\zeta$$

where

$$\hat{U}(x_\mu, x_3; x_3') = \pm H(\mp[x_3' - x_3]) \left\{ \prod_{\zeta=x_3'}^{x_3} \exp[-s\hat{\Gamma}(x_\mu, \zeta) d\zeta] \right\} \hat{u}(x_\mu, x_3')$$

The thin slab propagator

- ▶ The thin slab propagator
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then ,for a thin slab (Δx_3 , sufficiently small), we have the propagator

$$\hat{p}(\alpha_\mu, x_3; \alpha_\nu, x_3') \simeq \int \exp[is(\alpha_\sigma - \alpha_\sigma')x_\alpha] \exp[-s\hat{\gamma}(x_\mu, x_3 - \frac{1}{2}\Delta x_3, \alpha_\nu, s)\Delta x_3] dx_1 dx_2$$

The approximation method

The approximation method

Consider the 3-D acoustic one-way equation

$$\frac{\partial}{\partial z} U(x, y, z, \omega) = i \sqrt{\frac{\omega^2}{c^2} + \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}} U(x, y, z, \omega)$$

The approximation method

for which, the thin slab propagator is

$$g(x, y, z; x', y', z') \simeq \frac{1}{4\pi^2} \int \exp \left[i \sqrt{\frac{\omega^2}{c^2(x, y, z'')} - (k_x^2 + k_y^2)} \Delta z \right] \cdot \exp[i(k_x(x-x') + k_y(y-y'))]$$

where $z'' = z' + \frac{1}{2} \Delta z$

The approximation method

Let us take the kernel

$$\mathcal{A}(u, k) = \exp \left[i \sqrt{u^2 - k^2} dz \right]$$
$$u = \frac{\omega}{c(x, y, z)} \quad , \quad k^2 = k_x^2 + k_y^2$$

The approximation method

- ▶ the optimization approximation method is
- ▶ we get

The approximation method

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To find functions $\phi(u)$, $\psi(k)$ and a complex number λ , such that

$$\|\mathcal{A}(u, k) - \lambda\phi(u)\psi(k)^*\|_{L^2} = \min_{\bar{\phi}, \bar{\psi}, \bar{\lambda}} \|\mathcal{A}(u, k) - \bar{\lambda}\bar{\phi}(u)\bar{\psi}(k)^*\|_{L^2}$$

where

$$\begin{aligned}\bar{\phi} &\in \{\bar{\phi}(u) \in L^2[a, b], \|\bar{\phi}(u)\| = 1\} \\ \bar{\psi} &\in \{\bar{\psi}(u) \in L^2[c, d], \|\bar{\psi}(u)\| = 1\}\end{aligned}$$

$$\begin{aligned}\int_c^d \mathcal{A}(u, k)\psi(k)dk &= \lambda\phi(u) \\ \int_a^b \mathcal{A}(u, k)^*\phi(u)du &= \lambda^*\psi(k)\end{aligned}$$

The approximation method

- ▶ the optimization approximation method is

Which is transformed into an independent self-adjoint system of integral equations, given by

$$\int_a^b \int_c^d \mathcal{A}(u, k) \mathcal{A}(\bar{u}, k)^* \phi(\bar{u}) d\bar{u} dk = |\lambda|^2 \phi(u)$$
$$\int_a^b \int_c^d \mathcal{A}(u, k)^* \mathcal{A}(u, \bar{k}) \psi(\bar{k}) du d\bar{k} = |\lambda|^2 \psi(k)$$

The approximation method

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this integrals can be approximated by

$$\sum_{i=1}^m \left[\sum_{j=1}^n \Delta u \Delta k \mathcal{A}(u_l, k_j) \mathcal{A}(u_i, k_j)^* \right] \phi_i$$

The approximation method

- ▶ the optimization approximation method is

then, we have

$$F\phi = |\lambda|^2\phi$$

$G\psi = |\lambda|^2\psi$ for the second equation.

The approximation method

- ▶ the optimization approximation method is

Clearly

$$F = \mathcal{A}\mathcal{A}^H$$

$$G = \mathcal{A}^H\mathcal{A}$$

The approximation method

- ▶ the optimization approximation method is

and we have the approximation

$$\mathcal{A}(u, k) \simeq \lambda_1 \phi_1(u) \psi_1(k)^*$$

The approximation method

- ▶ the optimization approximation method is

which can be optimized as

$$\mathcal{A}(u, k) \simeq \sum_{l=1}^s \lambda_l \phi_l(u) \psi_l(k)^*$$

The approximation method

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$$g(x, y, z; x', y', z') = \frac{1}{4\pi^2} \sum_{l=1}^s \lambda_l \phi_l(u) \int \psi_l(k)^* \exp[i(k_x(x-x') + k_y(y-y'))] dk_x dk_y$$

Some references

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